

# INTRODUCTION TO MODERN COSMOLOGY

- (~ Dodelson Chapter 1)

Cosmology is really about answering the 3 oldest questions:

- Who are we?  $\longrightarrow$  What is the Universe made of / how did the elements form?
- Where do we come from?  $\longrightarrow$  How was the Universe born?
- Where are we going?  $\longrightarrow$  How did the Universe evolve, what is its fate?

Modern cosmology: quantitative, data-driven answers to these questions. Huge amount of data ( $\sim 10^7$  galaxies,  $10^8$  pixels) to test against predictions. Data-driven science, no longer in the realm of "philosophy".

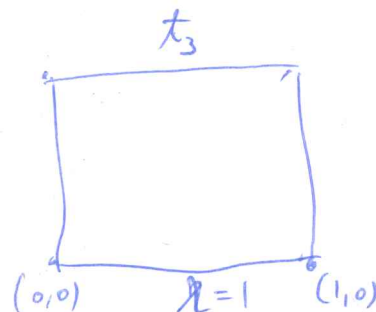
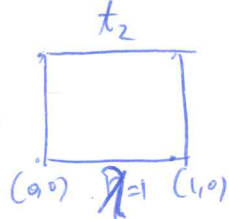
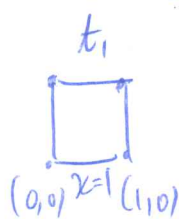
## The expanding Universe

Universe is expanding  $\longrightarrow$  distances between galaxies increase

Scale factor  $a$  Today  $a_0 = 1$ , in the past  $a < 1$   
"0": today

Physical distances  $\propto a$   
Comoving distances: fixed

Comoving grid

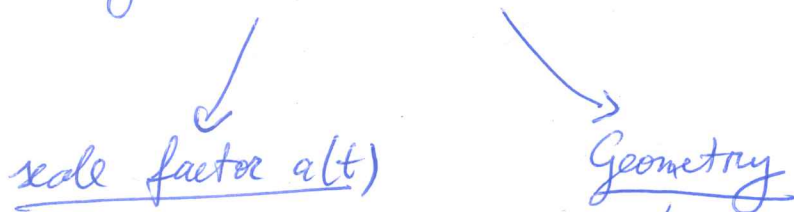


Comoving distances:

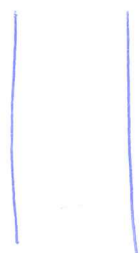
$t_1$	$t_2$	$t_3$
$\chi=1$	$\chi=1$	$\chi=1$
$a(t_1)\chi$	$a(t_2)\chi$	$a(t_3)\chi$

Physical distances

Smooth, expanding Universe characterized by ~~se~~



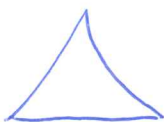
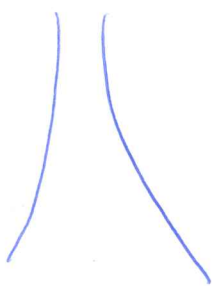
Flat



$180^\circ$

no curvature

Open



$<180^\circ$

negative curvature

Closed



$>180^\circ$

positive curvature

free paths of particles starting parallel to each other

GR: Geometry  $\equiv$  Energy

Flat Universe requires very special value for energy

density today:  $\rho_{crit} \sim 10^{-29} \frac{g}{cm^3}$

$\rho > \rho_{crit} \rightarrow$  closed

$\rho < \rho_{crit} \rightarrow$  open

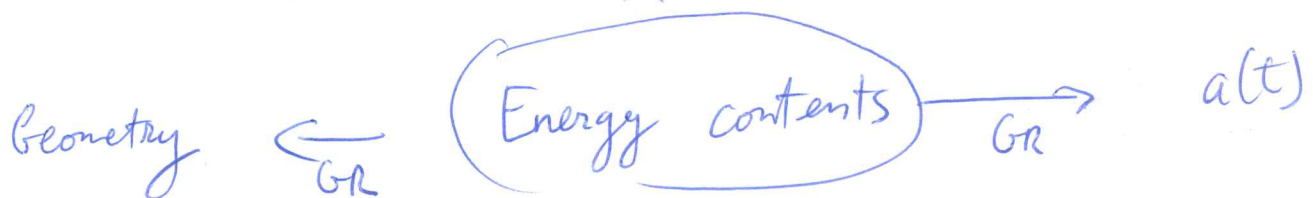
GR tells us not only how energy and geometry of the Universe are related, but in turn how  $a$  evolves in time

$a(t)$  depends on the dominant component of energy

Radiation ( $\gamma$ , massless  $\nu$ , etc.  $P = \frac{1}{3}\rho$ )  $\rightarrow a \propto t^{1/2}$

~~Matter~~  
Non-relativistic matter ("dust",  $P = 0$ )  $\rightarrow a \propto t^{2/3}$

Dark energy? ( $P = -\rho$ )  $\rightarrow a \propto e^t$



Hubble rate (logarithmic derivative of scale factor)

$$\boxed{H(t) \equiv \frac{1}{a} \frac{da}{dt}} = \frac{\dot{a}}{a}$$

Radiation domination:  $a \propto t^{1/2} \rightarrow H = \frac{1}{2t}$

Matter domination:  $a \propto t^{2/3} \rightarrow H = \frac{2}{3t}$

Dark energy domination:  $a \propto e^t \rightarrow H = \text{const}$  ( $a \propto e^{Ht}$ )

$H_0$  to powerful test of cosmology!

$\uparrow$  Age of Universe today

First Friedmann equation gives evolution of scale factor

(you'll see it again later, no need to worry now)

$$\boxed{H^2(t) = \frac{1}{a^2} \left( \frac{da}{dt} \right)^2 = \frac{8\pi G}{3} \left[ \rho(t) + \frac{\rho_{cr} - \rho_0}{a^2(t)} \right]} \quad \rho_{cr} \equiv \frac{3H_0^2}{8\pi G}$$

So the question of how  $a(t)$  evolves boils down to figuring out / providing a prescription for  $\rho(t)$

Non-relativistic matter ("dust"): "baryons", dark matter (DM)

$$E \propto m \quad \rho \propto \frac{E}{V} \propto mn \propto a^{-3} \quad \text{since } V \propto a^3$$

$$\rho_m \propto a^{-3}$$

Radiation

$$E \propto \frac{1}{\lambda} \propto \frac{1}{a} \quad \rightarrow \quad T(t) = \frac{T_0}{a(t)}$$

$\swarrow$   $k_B T_0$        $\swarrow$   $\frac{hc}{k_B T_0}$

Temperature of a thermal bath of radiation (e.g.  $\gamma$ )

$$\rho \propto \frac{E}{V} \propto \frac{1}{a} \frac{1}{a^3} \propto a^{-4} \quad [\rho \propto T^4 \text{ Stefan-Boltzmann law}]$$

$$\rho_r \propto a^{-4}$$

Dark energy

Cosmological constant  $\Lambda$ :  $\rho_\Lambda \propto \text{const}$  even as space expands

Example of getting  $a(t)$

Consider flat universe

$$H^2(t) = \frac{8\pi G}{3} \rho(t) \rightarrow \text{Radiation } \frac{1}{a^2} \left(\frac{da}{dt}\right)^2 \propto a^{-4}$$

$$\rightarrow \frac{1}{a} \frac{da}{dt} \propto a^{-2} \rightarrow \frac{da}{dt} \propto a^{-1} \rightarrow \int da a \propto \int dt \rightarrow t \propto a^2 \rightarrow a \propto t^{1/2}$$

$$\text{Matter } \frac{1}{a^2} \left(\frac{da}{dt}\right)^2 \propto a^{-3} \rightarrow \frac{1}{a} \frac{da}{dt} \propto a^{-3/2} \rightarrow \frac{da}{dt} \propto a^{-1/2} \rightarrow \int da a^{1/2} \propto \int dt$$

$$\rightarrow t \propto a^{3/2} \rightarrow a \propto t^{2/3}$$

Dark energy (cosmological constant)

$$\frac{1}{a} \left( \frac{da}{dt} \right)^2 = \text{const} = H_0^2 \quad \frac{1}{a} \frac{da}{dt} = H_0 \quad \int \frac{da}{a} = H_0 \int dt$$

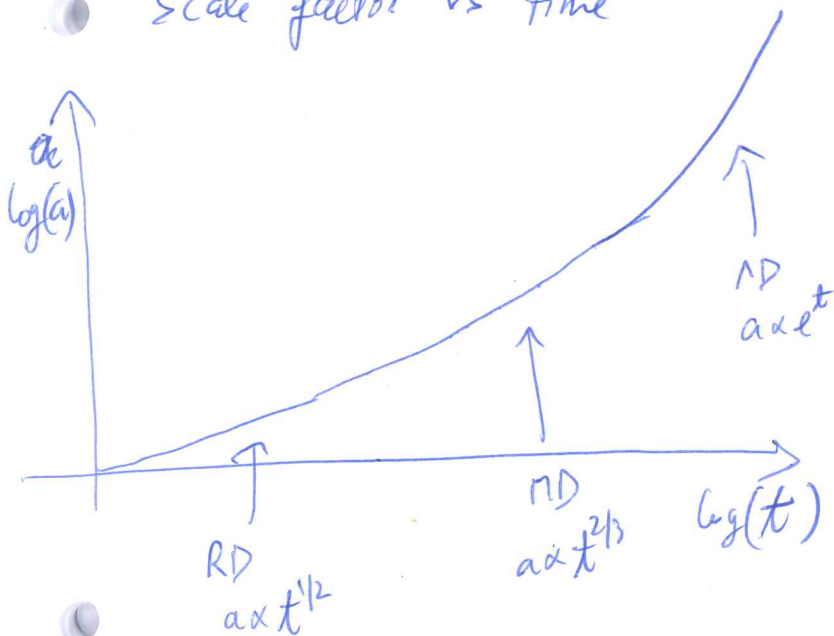
$$\ln a \propto H_0 t \rightarrow a \propto e^{H_0 t}$$

Recap

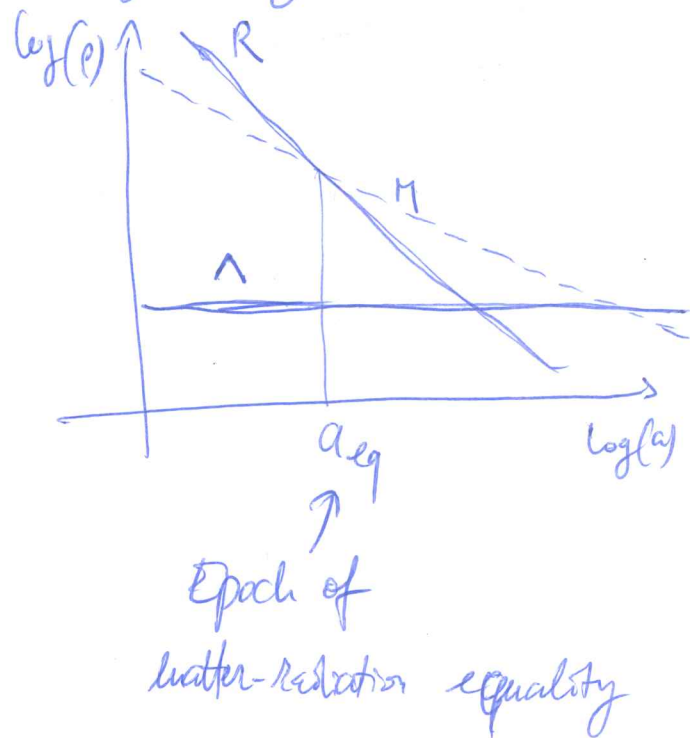
$$\rho_r \propto a^{-4} \rightarrow \rho_m \propto a^{-3} \rightarrow \rho_\Lambda \propto \text{const}$$

[curvature effectively behaves as  $\rho_k \propto a^{-2}$ ]

Scale factor vs time



Energy density vs scale factor



Some numbers

$$H_0 \sim 70 \frac{\text{km}}{\text{s} \cdot \text{Mpc}} \equiv 100 h \frac{\text{km}}{\text{s} \cdot \text{Mpc}}$$

$$1 \text{ Mpc} = 10^6 \text{ pc} \sim 3.1 \times 10^{24} \text{ cm}$$

$$h \sim 0.7 \quad (H_0 \sim 70 \frac{\text{km}}{\text{s} \cdot \text{Mpc}})$$

$$t_0 \sim 13 \text{ Gyr}$$

Easy to show  $H_0 t_0 > \frac{2}{3}$   
 $\downarrow$   
 $H_0 t_0$

# Redshift and Hubble's law

Expanding universe  $\rightarrow$  galaxies receding  $\rightarrow$  wavelength of emitted light stretched

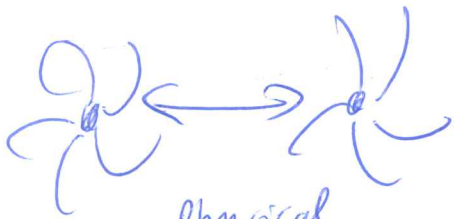
Stretching factor:  $H(z)$   $\leftarrow$  redshift

$$1+z \equiv \frac{\lambda_{obs}}{\lambda_{emit}} = \frac{1}{a} > 1$$

low redshift: Doppler formula  $z \approx \frac{v}{c}$

so  $\frac{\lambda_{obs}}{\lambda_{emit}} \approx 1 + \frac{v}{c}$  infer  
 $\uparrow$   
 measure

Hubble (1929): distant galaxies receding from us, recession velocity increases linearly with distance



physical distance  $d = aX$

assume galaxies fixed on the comoving grid, no peculiar velocity (only receding due to expansion)

$$v = \dot{d} = \left( \dot{a} X \right) = \dot{a} X + a \dot{X} =$$

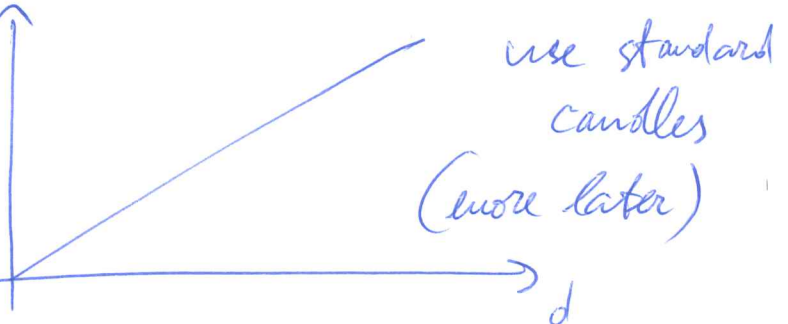
$$= \dot{a} X = \dot{a} \frac{d}{a} = \left( \frac{\dot{a}}{a} \right) d = H d$$

since  $H = \frac{1}{a} \frac{da}{dt}$

so  $v = Hd \propto d$

Evidence that  $\rightarrow$

Universe is expanding !!!



Beyond  $z \ll 1$ ,  $v = Hd$  no longer true

(later you will look at luminosity distance)

Hubble diagram:  $v$  vs  $d$  or more generally  $d$  vs  $z$

$z$ : easy to get, look at spectral lines

$d$ : hard, need objects with intrinsic fixed property (e.g. luminosity)

Type Ia Supernovae (SNe Ia) standard candles up to  $z \sim 3$

$d(z)$  relation (more later) depends on energy content

and cosmological model

### Equilibrium and beyond

$$\rho_r \propto a^{-4}$$

$$\rho_m \propto a^{-3}$$

Energy densities much higher in the past  $\Rightarrow$  hot, dense plasma initially in equilibrium

$\Rightarrow$  rates for particle interactions also much higher  $\Gamma$

$$\Gamma(z) \stackrel{?}{\gtrless} H(z)$$

$\Gamma_a(z) > H(z) \Rightarrow$  particle "a" in equilibrium

time

$\Gamma_a(z) < H(z) \Rightarrow$  particle "a" falls out of equilibrium or decouples

Early Universe very simple environment: smooth, hot, most

of its constituents in equilibrium

As it cools down particles start falling out of equilibrium ("freeze-out" / "decoupling")

$T \gtrsim \frac{1 \text{ MeV}}{k_B}$  : no neutral atoms, no bound nuclei

$\gamma$  immediately ionizes any atom/nucleus

$T \lesssim$  Binding energy  $\rightarrow$  lighter elements can form

Nuclear energies  $\sim$  MeV  $\rightarrow$  light nuclei form ( $D, {}^3\text{He}, {}^7\text{Li}, {}^4\text{He}$ )

Big Bang Nucleosynthesis (BBN)  $\leftarrow$  ( $t \sim 3 \text{ min}$ )  $\approx 10^{10}$

Can easily compute abundances of light elements just by knowing early Universe conditions + nuclear cross-sections

Light element abundances <sup>especially D</sup> depend on density of visible matter

(protons, neutrons, electrons  $\rightarrow$  baryons)  $\rho_b \propto a^{-3}$

need non-baryonic dark matter!  $\Leftarrow \rho_{b,0} \sim 5\% \rho_{\text{crit}}$

CMB: Cosmic Microwave Background

$T \sim \text{eV}, z \sim 1100$  ( $t \sim 300,000$  yrs)

$\rho_{\text{photon}}(z) < H(z)$   $\rightarrow$  photons can start propagating freely through space

Compton Scattering

$\downarrow$   
CMB: earliest "picture" of the Universe!

HUGE amount  $\leftarrow$  of information about the early Universe



CMB photons have blackbody spectrum

Specific intensity  $\rightarrow I_\nu = \frac{4\pi h \nu^3 / c^2}{\exp[2.82144 k_B T / h \nu] - 1}$

$\left[ \frac{\text{Energy}}{\text{Area} \cdot \text{time} \cdot \text{solid angle} \cdot \text{frequency}} \right]$

CMB temperature  $T \sim 3K$

CMB almost perfect  $T \sim 3K$  blackbody

Tiny ( $\sim 10^{-5}$ ) anisotropies in the sky encode crucial information

inhomogeneities  $\xrightarrow{\text{gravity}}$  structure

Dark matter, dark energy, and inflation

Hot Big Bang model pillars:

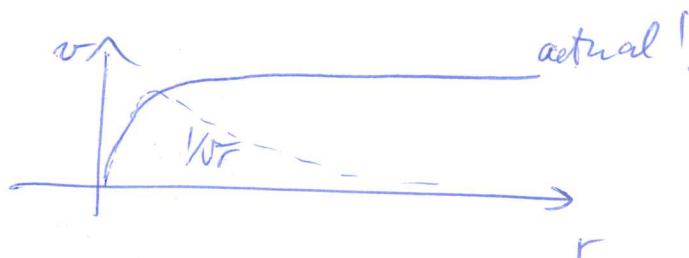
- Hubble's ~~law~~ diagram
- BBN
- CMB

But visible matter ("baryons") cannot be the end of the story

Need some non-luminous dark matter  $\left\{ \begin{array}{l} \text{BBN} \\ \text{CMB} \\ \text{rotation curves} \end{array} \right.$

Rotation curves of galaxies

$$\frac{1}{2} m v^2 = \frac{G M m}{r} \rightarrow v \propto \frac{1}{\sqrt{r}}$$



What is DM?

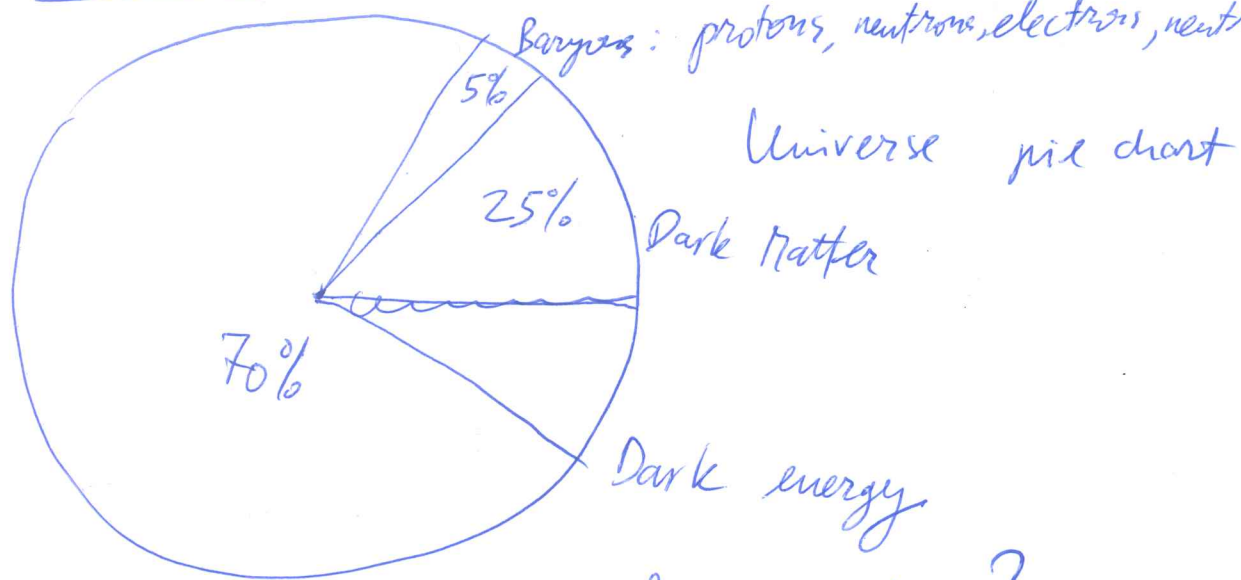
We don't know, but could be new elementary particle which once in equilibrium, then decoupled when  $T \sim GeV$

Also need dark energy to explain why current expansion is accelerated (we will see with just radiation and/or visible or dark matter, expansion is decelerated)

# We see lots of structure today: grew under gravitational instability starting from tiny seed perturbations which we see as anisotropies in the CMB

Who set the seeds for these perturbations?

Inflation! (also solves several other problems)  $t \sim 10^{-35}$



How to learn about all this from observations?

